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The structures and stability of silylenoids RBrSiLi₂ (R=CH₃, C(SiH₃)₃)

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Abstract The structures and stabilities of RBrSiLi₂ (R=CH₃, C(SiH₃)₃) have been studied using *ab initio* and DFT methods. CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ have three possible structures, the p-complex, the plain, and the folded structures. The plain and the folded structures are different from those of known structures of silylenoids. The energy of the plain structure is the lowest and nearly equals to that of the folded structure. The plain and the folded structures, which can isomerize into each other, are the most stable and possibly detected ones in chemical reactions. The ensertial of the insertion reactions with Me₃SiCl is the same. The insertion barriers are in the order of H₂SiLiBr > C(SiH₃)₃BrSiLi₂ > CH₃BrSiLi₂. The C(SiH₃)₃ group makes the insertion of C(SiH₃)₃BrSiLi₂ more difficult.

Keywords $Ab initio \cdot DFT \cdot Insertion reactions \cdot Silylenoids \cdot Theoretical Study$

Introduction

Silylenoids, R₂SiMX (X=halogen, M=alkali metal), are important intermediates in silicon hybrid and organosilicon chemistry [1, 2]. Previous research works have shown that silylenoids can be regarded as a silylene complex, in which a

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Environment Research Institute, Shandong University, Jinan 250100, People's Republic of China leaving group X and a metal atom M are bound to the same silicon atom. Clark et al. [3] have carried out the first theoretical study on the simplest silylenoid H₂SiLiF. During recent years, the silylenoids such as R₂SiMX (R=CH₃, C₂H₅, NH₂, OH, H; X=F, Cl, Br; M=Li, Na, K) have been systematically studied [4–10]. Their structures, stabilities, and the factors of influcencing their existence have been investigated in detail. These research works show that though silylenoids are very reactive, they are detectable and their stable structures are not "classical" tetrahedral structures (S1), but are p-complex (S2) and three-membered ring (S3) structures. These results will provide useful and important theoretical information for the understanding of their reaction mechanisms.



As a kind of active species, the preparation of silylenoids is difficult. In 1995, Tamao et al. experimentally detected the existence of silylenoid Ph₂SiLi(OBu-t) for the first time [11]. Most recently, a great breakthrough in the research of silylenoids has been made. Lee et al. reported the syntheses of stable halosilylenoids (Tsi)X₂SiLi (Tsi=C(SiMe₃)₃, X=Br, Cl) at room temperature [12].

Recently, Myong Euy Lee reported his experimental research on the reaction of silylenoids (Me₃Si)₃CBrSiLi₂ in Shandong University. (Me₃Si)₃CBrSiLi₂ is another kind of silylenoid. At present, the structures and chemical properties of this silylenoid have not been studied. According to the components of the silylenoid, it can be expected that (Me₃Si)₃CBrSiLi₂ certainly has its own specific properties.

What are the structures and chemical properties of (Me₃Si)₃CBrSiLi₂? To get insight into these questions and provide useful theoretical information for the deep studying of analogues

of (Me₃Si)₃CBrSiLi₂, CH₃BrSiLi₂, and C(SiH₃)₃BrSiLi₂ are nvestigated and analyzed by density functional theory in the present paper.

Theoretical methods

The structures of reactants, transition states, and products were fully optimized at the B3LYP/6-31G(d) level [13-16]. The corresponding harmonic vibrational frequency calculations were carried out in order to characterize all stationary points as either local minima (no imaginary frequencies) or transition states (one imaginary frequency). Each transition state was verified to connect the designated reactants and products by performing an intrinsic reaction coordinate (IRC) analysis (The step size is 0.1 amu^{1/2}-Bohr. The number of steps is 70.) [17, 18]. Based on the optimized geometries, energies were obtained and natural bond orbital (NBO) [19-21] analyses were then used to study the nature of different interactions between atoms and groups. The solvent effects, which were simulated using the self-consistent reaction field (SCRF) method with Tomasi's polarized continuum model (PCM) [22–31], were investigated at the same level. The structures of CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ are also optimized at the MP2 (full)/6-31G(d) level [32] and G3(MP2) [33] theory was used for calculations of molecular energies. The Gaussian 03 [34] series of programs were employed in all calculations.

Results and discussion

Isomer structures of silylenoids

The structures of CH₃BrSiLi₂, C(SiH₃)₃BrSiLi₂ in vacuum are located at the B3LYP/6-31G(d) and the MP2(full)/6-31G(d) levels firstly. The PCM method is used to determine the structures in THF at the B3LYP/6-31G(d) level. Based on the MP2(full)/6-31G(d) geometries, energies were obtained using the G3(MP2) theory. Standard orientations and thermodynamic results of stationary points, imaginary frequencies of transition states are listed in Supporting information. The B3LYP/6-31G(d), MP2(full)/6-31G(d), and G3(MP2) energies are shown in Tables 1 and 2. Figures 1 and 2 indicate the calculated structures at the B3LYP/6-31G(d) level. For convenience, the central Si atom in silylenoids is marked as Si¹.

Structures of silylenoids in vacuum

Figure 1 shows the structures of $CH_3BrSiLi_2$ and $C(SiH_3)_3BrSiLi_2$ in vacuum at the B3LYP/6-31G(d) level. Unless otherwise specified, the following results are based on the calculations at the B3LYP/6-31G(d) level.

CH₃BrSiLi₂ Two structures of CH₃BrSiLi₂ are located. One is the p-complex structure 1a, the other is the plain structure 1b. 1a can be regarded as the result of the interaction between the negative Br of LiBr and the empty p-orbital of Si in CH₃LiSi. The strong interaction between Br and Si makes 1a very stable. The energy of **1a** is 128.2 kJ mol⁻¹ lower than the sum of the separated LiBr and CH₃LiSi molecules. 1a is similar to that of the p-complex structure of H₂SiLiBr [3–9]. In the plain structure **1b**, the Si^1 atom is sp^2 hybridized. Three sp² hybridized orbitals are connected with C and two Li atoms, respectively. Si¹, C, Br, and two Li atoms are all in the same plain. Their natural charges are -1.213 (C), -0.204(Si¹), -0.825 (Br), 0.740 (Li). NBO analyses indicate that the main parts of HOMO and LUMO in **1b** are situated on the Si¹ atom. The main HOMO is localized on the p-orbital of the Si¹-Li-Br-Li plain (occupied number: 0.502), and the main LUMO is on the p-orbital vertical to the plain (occupied number: 0.506). **1b** is 66.4 kJ mol⁻¹ lower in energy than 1a. To the best of our knowledge, the structure of 1b is different from any known structures of silvlenoids.

 $C(SiH_3)_3BrSiLi_2$ Calculation results show that $C(SiH_3)_3BrSiLi_2$ has three isomers, the p-complex structure (2a), the plain structure (2b), and the folded structure (2c). Similar to 1a, 2a can be regarded as the complex of LiBr and C(SiH₂)₃LiSi. The energy of 2a is 125.1 kJ mol⁻¹ lower than the sum of the separated LiBr and C(SiH₃)₃LiSi molecules. The structure 2b is similar to 1b. Relative to the methyl in 1b, 2b bears the stronger electrondonating group, $C(SiH_3)_3$. The natural charges of Si^1 (-0.439) in **2b** is lower than that (-0.204) in **1b**. Similar to **1b**, the main parts of HOMO and LUMO of **2b** are situated on the atom Si¹. One is the p-orbital of the Si-Li-Br-Li plain (occupied number: 0.485), the other is the p-orbital vertical to the plain (occupied number: 0.637). The energy of **2b** is 64.4 kJ mol⁻¹ lower than that of **2a**. In 2c, Si¹, Br and two Li atoms are in one plain and the bond of Si¹-C is almost vertical to the plain. The main structure difference between **2b** and **2c** is the position of the $C(SiH_3)_3$ group. The Si¹ atom is also sp^2 hybridized in **2c**. Two sp^2 hybridized orbital of Si¹ are connected with two Li atoms, respectively, and the porbital of Si¹ interacts with the C atom of the C(SiH₃)₃ group. 2c isomerizes into 2b via 2cbts with only a small barrier of 7.0 kJ mol^{-1} ; 22.9 kJ mol^{-1} is needed for **2b** to isomerize into **2c**. So, it is shown that structure 2b is more stable than 2c. The energy of **2c** is 15.9 kJ mol⁻¹ higher than that of **2b**. The main parts of HOMO and LUMO in 2c are also localized on the central Si¹. The occupied number of the p-orbital on the Si-Li-Br-Li plain (HOMO) is 0.486, and the occupied number of the p-orbital vertical to the plain (LUMO) is 0.441.

In summary, **1a** and **2a** can be regarded as the complexes of silylenes and LiBr, resulted from the interaction between the empty p-orbital of the Si¹ atom and the negative Br atom. In **1b**, **2b**, and **2c**, the Si¹ atom, which is sp² hybridized, is the active center. The

Table 1 Total energies (a.u.) and relative energies $(kJ mol^{-1}, in parentheses)$ for CH_3SiLi_2Br,C $(SiH_3)_3SiLi_2Br$ in vacuum and in THF at the B3LYP/6-31G(d) level

Molecules	Е	ZPE	E + ZPE
In vacuum			
1a	-2916.21319(0.0)	0.04050	-2916.17270(0.0)
1b	-2916.23829(-65.9)	0.04029	-2916.19800(-66.4)
$CH_3LiSi + LiBr$	-2916.16140(135.9)	0.03755	-2916.12385(128.2)
2a	-3788.31767(0.0)	0.08796	-3788.22971(0.0)
2b	-3788.34142(-62.3)	0.08719	-3788.25423(-64.4)
2cbts	-3788.33255(-39.1)	0.08704	-3788.24552(-41.5)
2c	-3788.33607(-48.3)	0.08790	-3788.24818(-48.5)
C(SiH ₃) ₃ LiSi + LiBr	-3788.26759(131.5)	0.08554	-3788.18206(125.1)
In THF			
1a	-2916.24406(0.0)	0.03989	-2916.20416(0.0)
1b	-2916.25718(-34.4)	0.03937	-2916.21781(-35.8)
1c	-2916.25423(-26.7)	0.03930	-2916.21494(-28.3)
2a	-3788.34478(0.0)	0.08684	-3788.25794(0.0)
2b	-3788.35998(-39.9)	0.08598	-3788.27400(-42.2)
2c	-3788.35736(-33.0)	0.08682	-3788.27054(-33.1)

energies are in the orders of 1a > 1b,2a > 2c > 2b. The energy difference between 1a and 1b (66.4 kJ mol⁻¹) is almost as large as that between 2a and 2b (64.4 kJ mol⁻¹). However, the energy difference (15.9 kJ mol⁻¹) between 2b and 2c is much smaller. Both the calculated structures and energies of CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ indicate that the plain structure is close to the folded structure and has large difference from the p-complex structure. **1b** and **2b** are the most stable and possibly detectable ones in chemical reactions.

The calculated structures of CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ at the MP2(full)/6-31G(d) level are similar to those calculated at the B3LYP/6-31G(d) level. The MP2(full)/6-31G(d) level and G3(MP2) energies also indicate that the order of 1a > 1b, 2a > 2c > 2b.

Table 2 Total energies (a.u.) and relative energies $(kJ mol^{-1}, in parentheses)$ for $CH_3SiLi_2Br, C(SiH_3)_3SiLi_2Br$ in vacuum

Molecules	MP2(full)/6-31G(d)	G3 (MP2)
1a	-2913.71678 (0.0)	-2913.76479(0.0)
1b	-2913.73958(-58.9)	-2913.78751(-59.5)
CH3LiSi + LiBr	-2913.65884(152.1)	-2913.70895(146.6)
2a	-3784.21165(0.0)	-3784.32385(0.0)
2b	-3784.23111(-51.2)	-3784.34265(-49.4)
2cbts	-3784.22756(-41.8)	-3784.33898(-39.7)
2c	-3784.22933(-46.5)	-3784.34136(-46.0)
C(SiH ₃) ₃ LiSi + LiBr	-3784.15118(158.7)	-3784.26772(147.2)

Structures of silylenoids in THF

The PCM model at the B3LYP/6-31G(d) level is used to calculate the structures of CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ (see Supporting information). The results can be summarized as follows:

- Three structures, the p-complex, the plain, and the folded structures of CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ in THF are located. These structures are similar to their corresponding ones in vacuum.
- (2) Different from that in vacuum, the folded structure 1c of CH₃BrSiLi₂ is located in THF. Just like 2c resembling 2b, 1c is close to 1b both in their energies (see Table 1) and in their structures.
- (3) The thermal stabilities of the structures are lower in THF than in vacuum. The energies in THF are in the same order of 1a > 1c > 1b, 2a > 2c > 2b as those in vacuum.

Insertion reactions with Me₃SiCl

To gain a further understanding about the stability of these silylenoids, the insertion reactions of $CH_3BrSiLi_2$, $C(SiH_3)_3BrSiLi_2$ with Me_3SiCl have been investigated at the B3LYP/6-31G(d) level. The insertion of the simplest model of silylenoid, $H_2SiLiBr$, is also studied for comparison. The optimized geometries are shown in Figs. 2 and 3 and Supporting information. The total energies together with the zero-point energies (ZPEs) and relative energies (relative to the corresponding reactants) of all stationary points are

Fig. 1 The B3LYP/6-31G(d) geometries (in Å and (°)) for CH_3SiLi_2Br , $C(SiH_3)_3SiLi_2Br$, and the transition state 2cbts in vacuum. The central Si atom in silylenoids is marked as Si¹



described by Table 3. For convenience, the Si atom in Me_3SiCl is marked as Si^2 .

Insertion reaction of H₂SiLiBr with Me₃SiCl

The previous calculations [3-9] show that silylenoid H₂SiLiBr has four equilibrium isomers, the three-membered ring, the p-complex, the σ -complex (linear), and the "classical" tetrahedral structures. The three-membered ring structure is the most stable and possibly detectable one in chemical reactions. So the three-membered ring structure of silylenoid H₂SiLiBr is adopted in this study.

Transition state Two electron donation effects contribute to the proceeding of the insertion reaction. One is the donation of the electrons of Cl into the p-orbital on the Si¹ atom. The other is the donation of the σ electrons on the Si¹ atom to the positive Me₃Si² group. The electron donations make the formation of the transition state **INts**. The insertion barrier is 64.1 kJ mol⁻¹.

Insertion product After getting over the transition state **INts**, **INpro** are gradually formed with the LiBr moiety leaving from the Si¹ atom. In fact, **INpro** is a complex of **silane1** and LiBr. The energy of **INpro** is 34.9 kJ mol⁻¹ lower than the sum of **silane1** and LiBr.

Insertion reaction of CH₃BrSiLi₂ with Me₃SiCl

Precursor complex Different from the process of H₂SiLiBr insertion, there is the precursor complex **11Npre** formed in the insertion reaction of CH₃BrSiLi₂ with Me₃SiCl. In the beginning of the insertion, the negative Cl of Me₃SiCl approaches the p-orbital of Si¹ from one side of the LiSi¹Li plain, and the precursor complex **11Npre** forms. In **11Npre**, the Cl···Si¹ distance is 2.921 Å, and the Cl...Si² distance (2.140 Å) is only slightly different from that of Me₃SiCl moieties in **11Npre** is very weak. The energy of **11Npre** is only 19.7 kJ mol⁻¹ lower than the sum of CH₃BrSiLi₂ and Me₃SiCl.



1INpro

silylenoid1

Transition state As the Cl^{...}Si¹ interaction becomes strong, the positive Me₃Si group of Me₃SiCl simultaneously interacts with

the negative Si^1 atom of $CH_3BrSiLi_2$. The insertion reaches the transition state **1INts**. In **1INts**, the Si^1 -Cl and Si^1 - Si^2 distances

Fig. 3 The B3LYP/6-31G(d) geometries (in Å and (°)) for some stationary points in the insertion reaction of C(SiH₃)₃SiLi₂Br with Me₃SiCl. The central Si atom in silylenoids and the Si atom in Me₃SiCl are marked as Si¹ and Si², respectively



Table 3 Total energies (a.u.) and relative energies $(kJ \text{ mol}^{-1}, in \text{ parentheses})$ for reactants, intermediates, transition states, and products of the insertion reactions at the B3LYP/6-31G(d) level

Molecules	Е	ZPE	E + ZPE
H ₂ SiLiBr + Me ₃ SiCl	-3739.50603(0.0)	0.13147	-3739.37456(0.0)
INts	-3739.48333(59.6)	0.13317	-3739.35016(64.1)
INpro	-3739.53878(-86.0)	0.13384	-3739.40494(-79.7)
Silane1+ LiBr	-3739.52417(-47.6)	0.13255	-3739.39162(-44.8)
1b + Me ₃ SiCl	-3785.76085(0.0)	0.15381	-3785.60704(0.0)
1INpre	-3785.76932(-22.2)	0.15478	-3785.61454(-19.7)
1 INts	-3785.75644(11.6)	0.15511	-3785.60133(15.0)
1INpro	-3785.85495(-247.0)	0.15532	-3785.69963(-243.0)
Silylenoid1+ LiBr	-3785.77787(-44.7)	0.15379	-3785.62407(-44.7)
$\mathbf{2b} + Me_3SiCl$	-4657.86398(0.0)	0.20071	-4657.66327(0.0)
2INts	-4657.84721(44.0)	0.20199	-4657.64522(47.4)
2INpro	-4657.94392(-209.8)	0.20271	-4657.74122(-252.0)
Silylenoid2 + LiBr	-4657.86803(-10.6)	0.20143	-4657.66659(-8.7)

are 0.395 and 1.300 Å shorter than those in **1INpre**, respectively. The breaking Cl-Si² distance is 2.243 Å. The energy of **1INts** is 15.0 kJ mol^{-1} relative to the sum of CH₃BrSiLi₂ and Me₃SiCl.

Insertion product As the reaction proceeds, the Si¹-Cl and Si¹-Si² bonds are formed with the breaking of the Cl-Si² bond. The insertion product silane **1INpro** formed. **1INpro** can be decomposed into **silylenoid1** and LiBr. The energy (198.3 kJ mol⁻¹) needed for the dissociation of **1INpro** is far less than that given out by the process from **1INts** to **1INpro** (258.0 kJ mol⁻¹). So there is enough energy in the reaction system to dissociate **1INpro** into **silylenoid1** and LiBr.

The structures of **silylenoid1** are similar to those of $H_2SiLiBr$. The same as with $H_2SiLiBr$, **silylenoid1** can react with Me₃SiCl.

Insertion reaction of C(SiH₃)₃BrSiLi₂ with Me₃SiCl

The insertion of $C(SiH_3)_3BrSiLi_2$ with Me₃SiCl is similar to the insertion of $CH_3BrSiLi_2$ with Me₃SiCl, but there is no precursor complex formed in the insertion process of $C(SiH_3)_3BrSiLi_2$. The relative energy of the transition state **2INts** is 47.4 kJ mol⁻¹. The insertion product **2INpro** can decompose into **silylenoid2** and LiBr. Just like H₂SiLiBr, **silylenoid2** can react with Me₃SiCl.

The above discussion shows that the insertion processes of H₂SiLiBr, CH₃BrSiLi₂, and C(SiH₃)₃BrSiLi₂ with Me₃SiCl are similar. Essential to these insertions are the p-orbital and electrons of Si¹ in silylenoids which interact with the negative Cl and positive Me₃Si group, respectively. In the view of energy, the insertion barriers are in the order of CH₃BrSiLi₂ < C(SiH₃)₃BrSiLi₂ < H₂SiLiBr. This indicates that silylenoids of the CH₃BrSiLi₂ analogues are easier to insert with Me₃SiCl than those of the H₂SiLiBr style. The big C(SiH₃)₃ group makes the insertion of C(SiH₃)₃BrSiLi₂ more difficult than that of CH₃BrSiLi₂. NBO analyses indicate that the occupancy of p-orbital of Si¹ in CH₃BrSiLi₂ (0.048) is much lower than that in C(SiH₃)₃BrSiLi₂ (0.285). Calculations also show that the natural charge of Si¹ in CH₃BrSiLi₂ (-0.204) is higher than in that C(SiH₃) ₃BrSiLi₂ (-0.439). The lower occupancy of p-orbital and the lower charge of Si¹ benefit to the insertion reaction of silylenoids. So the C(SiH₃)₃ group make the p-orbital accept electrons more difficult and the negative Si interact on positive groups easier. The combined effects result in the higher insertion barrier of C(SiH₃)₃BrSiLi₂.

Conclusions

In the present work, we have studied the structures of silylenoids CH₃BrSiLi₂, C(SiH₃)₃BrSiLi₂ and their insertion reactions with Me₃SiCl by *ab initio* and DFT theory.

- (1) The calculations indicate that CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂ have three possible structures: the p-complex, the plain, and the folded structures. The plain and the folded structures are different from those of known structures of silylenoids. The energy of the plain structure is the lowest and nearly equals to that of the folded structure. The plain and the folded structures are possibly the detected ones in chemical reactions of CH₃BrSiLi₂ and C(SiH₃)₃BrSiLi₂.
- (2) The insertions of H₂SiLiBr, CH₃BrSiLi₂ and C(SiH₃) ₃BrSiLi₂ into Si-Cl bond are similar. The essential of these insertion reactions is same. The insertion barriers are in the order of H₂SiLiBr > C(SiH₃)₃BrSiLi₂ > CH₃BrSiLi₂. The C(SiH₃)₃ group makes the insertion more difficult.

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References

- 1. Gilman H, Peterson DJ (1965) J Am Chem Soc 87:2389-2394
- 2. Nefedow OM, Manakow MN (1964) Angew Chem 76:270
- 3. Clark T, Schleyer PJ (1980) Organomet Chem 191:347-353
- 4. Feng DC, Feng SY, Deng CH (1995) Chem J Chin Univ 17: 1108–1111
- 5. Feng SY, Zhou YF, Feng DCJ (2003) Phys Chem A 107(20): 4116-4121
- 6. Qi Y, Feng D, Feng S (2008) J Mol Struct (Theochem) 856: 96–104
- 7. Xie J, Feng D, Feng SJ (2006) Organomet Chem 691(1-2):208-223
- 8. Xie J, Feng D, Feng S, Zhang J (2005) J Mol Struct (Theochem) 755(1–3):55–63
- 9. Feng S, Feng D, Li J (2000) Chem Phys Lett 316(1-2):146-150
- 10. Feng S, Deng C (1991) Chem Phys Lett 186(2–3):248–252
- Tamao K, Asahara M, Saeki T, Toshimitsu A (1999) Angew Chem Int Ed Engl 38:3316–3318
- Lee ME, Hyeon MC, Lim YM, Choi JK, Park CH, Jeong SE, Lee U (2004) Chem Eur J 10:377–381
- 13. Beck A (1993) J Chem Phys 98:5648–5652
- 14. Beck A (1988) Phys Rev A 38:3098–3100
- 15. Vosko S, Wilk L, Nusair M (1980) Can J Phys 58:1200-1211
- 16. Lee C, Yang W, Parr R (1988) Phys Rev B 37:785–789
- 17. Gonzalez C, Schlegel HB (1989) J Chem Phys 90:2154
- 18. Gonzalez C, Schlegel HBJ (1990) Phys Chem 94:5523-5527
- 19. Foster J, Weinhold FJ (1980) Am Chem Soc 102:7211–7218
- 20. Reed A, Weinhold FJ (1983) Chem Phys 78:4066-4073

- Weinhold F (1998) In: Schleyer PVR (Ed.) Encyclopedia of Computational Chemistry, pp. 1792–1811
- 22. Miertus S, Scrocco E, Tomasi J (1981) Chem Phys 55:117-129
- 23. Miertus S, Tomasi J (1982) Chem Phys 65:239-245
- 24. Cossi M, Barone V, Cammi R, Tomasi J (1996) Chem Phys Lett 255: 327–335
- 25. Cancès MT, Mennucci B, Tomasi JJ (1997) Chem Phys 107: 3032–3036
- 26. Barone V, Cossi M, Tomasi J (1997) J Chem Phys 107: 32103221
- 27. Cossi M, Barone V, Mennucci B, Tomasi J (1998) Chem Phys Lett 286:253–263
- 28. Barone V, Cossi M, Tomasi J (1998) J Comp Chem 19:404-417
- 29. Barone V, Cossi M (1998) J Phys Chem A 102:1995-2001
- 30. Mennucci B, Tomasi JJ (1997) Chem Phys 106:5151–5160
- Tomasi J, Mennucci B, Cancès E (1999) J Mol Struct (Theochem) 464:211–226
- 32. Lee C, Yang W, Parr R (1988) Phys Rev B 37:785–789
- Curtiss LA, Redfern PC, Raghavachari K, Rassolov V, Pople JAJ (1999) Chem Phys 110:4703–4705
- 34. Frisch MJ, Trucks GW, Schlegel HB, Scuseria GE, Robb MA, Cheeseman JR, Montgomery JA, Vreven NJ, Kudin TK, Burant JC, Millam JM, Iyengar SS, Tomasi JB, Mennucci V, Cossi BM, Scalmani GN, Rega G, Petersson A, Nakatsuji H, Hada M, Ehara M, Toyota K, Fukuda R, Hasegawa J, Ishida M, Nakajima T, Honda Y, Kitao O, Nakai H, Klene M, Li X, Knox JE, Hratchian HP, Cross JB, Bakken V, Adamo C, Jaramillo J, Gomperts R, Stratmann RE, Yazyev O, Austin AJ, Cammi R, Pomelli C, Ochterski JW, Ayala PY, Morokuma K, Voth GA, Salvador P, Dannenberg J, Zakrzewski VG. Dapprich S. Daniels AD. Strain MC. Farkas O. Malick DK. Rabuck AD, Raghavachari K, Foresman JB, Ortiz JV, Cui Q, Baboul AG, Clifford S, Cioslowski J, Stefanov BB, Liu G, Liashenko A, Piskorz P, Komaromi IR, Martin LD, Fox J, Keith T, AlLaham MA, Peng CY, Nanayakkara A, Challacombe M, Gill PM, Johnson WB, Chen W, Wong MW, Gonzalez C, Pople JA (2003) Gaussian 03, revision. Gaussian, Inc., Pittsburgh